

Now consider the case of isotropic pressure. To calculate the pressure, consider the force  $dF$  in the  $z$  direction of the particles on a piston (or the side of a box containing the particles) of small area  $dA$  that is perpendicular to the  $z$  direction (note: in this subsection  $z$  denotes a Cartesian coordinate, not redshift). The pressure is defined as

$$p = \frac{dF}{dA} = \frac{dq_z/dt}{dA} ,$$

where  $dq_z$  is the  $z$  component of momentum imparted to the piston, and we used the  $z$  component of Newton's law in the form that is also valid in special relativity. Particles with  $z$  momentum  $q_z$  contribute  $2q_z$  per particle to  $dq_z$  (with the factor of 2 due to elastic recoil). Particles with  $z$  velocity  $v_z$  can reach the piston during the time  $dt$  from as far away as  $dz = v_z dt$ . Thus,

$$dq_z = \int 2q_z f d^3q dA v_z dt .$$

Now, in general (in special relativity)  $\vec{q} = \gamma m \vec{v}$ , where  $m$  is the particle mass and  $\gamma = 1/\sqrt{1-v^2}$ . Also, isotropy implies:

$$\langle v_x^2 \rangle = \langle v_y^2 \rangle = \langle v_z^2 \rangle = \frac{1}{3} \langle v^2 \rangle ,$$

where the averaging here is over all the particles at a given position. Thus,  $\langle q_z v_z \rangle = \frac{1}{3} \langle q v \rangle$ , but we must then add a factor of 1/2 since in this averaging we must count only the half of particles headed towards the piston (rather than the opposite  $z$  direction). Thus,

$$dq_z = \frac{1}{3} \int q v f d^3q dA dt ,$$

so that

$$p = \frac{1}{3} \int q v f d^3q . \tag{2.29}$$