

Galaxies

Lecture 11

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<http://wise-obs.tau.ac.il/~barkana/galaxies.html>

Euler equation: Gas disk ($z = 0$)
(Include gravity and pressure)

$$\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} = -\frac{1}{\rho} \vec{\nabla} p - \vec{\nabla} \Phi$$

In cylindrical coordinates:

$$\frac{\partial v_R}{\partial t} + v_R \frac{\partial v_R}{\partial R} + \frac{v_\phi}{R} \frac{\partial v_R}{\partial \phi} - \frac{v_\phi^2}{R} = -\frac{\partial \Phi}{\partial R} - \frac{1}{\Sigma} \frac{\partial p}{\partial R}$$

$$\frac{\partial v_\phi}{\partial t} + v_R \frac{\partial v_\phi}{\partial R} + \frac{v_\phi}{R} \frac{\partial v_\phi}{\partial \phi} + \frac{v_\phi v_R}{R} = -\frac{1}{R} \frac{\partial \Phi}{\partial \phi} - \frac{1}{\Sigma R} \frac{\partial p}{\partial \phi}$$

Now assume a simple equation of state: $p = K\Sigma^\gamma$

The speed of sound: $c_s^2 = \frac{dp}{d\Sigma} = \gamma K \Sigma^{\gamma-1}$

Now simplify using specific enthalpy h :

$$\frac{\partial v_R}{\partial t} + v_R \frac{\partial v_R}{\partial R} + \frac{v_\phi}{R} \frac{\partial v_R}{\partial \phi} - \frac{v_\phi^2}{R} = - \frac{\partial}{\partial R} (\Phi + h)$$

$$\frac{\partial v_\phi}{\partial t} + v_R \frac{\partial v_\phi}{\partial R} + \frac{v_\phi}{R} \frac{\partial v_\phi}{\partial \phi} + \frac{v_\phi v_R}{R} = - \frac{1}{R} \frac{\partial}{\partial \phi} (\Phi + h)$$

Unperturbed equation: $\frac{v_{\phi,0}^2}{R} = \frac{d}{dR} (\Phi_0 + h_0)$

Perturbed variables:

$$\begin{aligned} v_R &= v_{R,1} & v_\phi &= v_{\phi,0} + v_{\phi,1} \\ h &= h_0 + h_1 & \Phi &= \Phi_0 + \Phi_1 \end{aligned}$$

1'st order equations (subtracted 0'th order):

$$\frac{\partial v_{R,1}}{\partial t} + \Omega \frac{\partial v_{R,1}}{\partial \phi} - 2\Omega v_{\phi,1} = - \frac{\partial}{\partial R} (\Phi_1 + h_1)$$

$$\frac{\partial v_{\phi,1}}{\partial t} + \left[\frac{d(\Omega R)}{dR} + \Omega \right] v_{R,1} + \Omega \frac{\partial v_{\phi,1}}{\partial \phi} = - \frac{1}{R} \frac{\partial}{\partial \phi} (\Phi_1 + h_1)$$

$$\frac{\partial v_{R,1}}{\partial t} + \Omega \frac{\partial v_{R,1}}{\partial \phi} - 2\Omega v_{\phi,1} = -\frac{\partial}{\partial R} (\Phi_1 + h_1)$$

$$\frac{\partial v_{\phi,1}}{\partial t} + \left[\frac{d(\Omega R)}{dR} + \Omega \right] v_{R,1} + \Omega \frac{\partial v_{\phi,1}}{\partial \phi} = -\frac{1}{R} \frac{\partial}{\partial \phi} (\Phi_1 + h_1)$$

Search for wave solutions:

$$v_{R,1} = \text{Re} \left[v_{R,a}(R) e^{i(m\phi - \omega t)} \right]$$

$$v_{\phi,1}; \Phi_1; h_1; \Sigma_1$$

• • • Substitute; Solve for velocity components =>

$$v_{R,a} = -\frac{i}{\Delta} \left[(m\Omega - \omega) \frac{d}{dR} (\Phi_a + h_a) + \frac{2m\Omega}{R} (\Phi_a + h_a) \right]$$

$$v_{\phi,a} = \frac{1}{\Delta} \left[-2B \frac{d}{dR} (\Phi_a + h_a) + \frac{m(m\Omega - \omega)}{R} (\Phi_a + h_a) \right]$$

$$v_{R,a} = -\frac{i}{\Delta} \left[(m\Omega - \omega) \frac{d}{dR} (\Phi_a + h_a) + \frac{2m\Omega}{R} (\Phi_a + h_a) \right]$$

$$v_{\phi,a} = \frac{1}{\Delta} \left[-2B \frac{d}{dR} (\Phi_a + h_a) + \frac{m(m\Omega - \omega)}{R} (\Phi_a + h_a) \right]$$

$$\Delta \equiv \kappa^2 - (m\Omega - \omega)^2$$

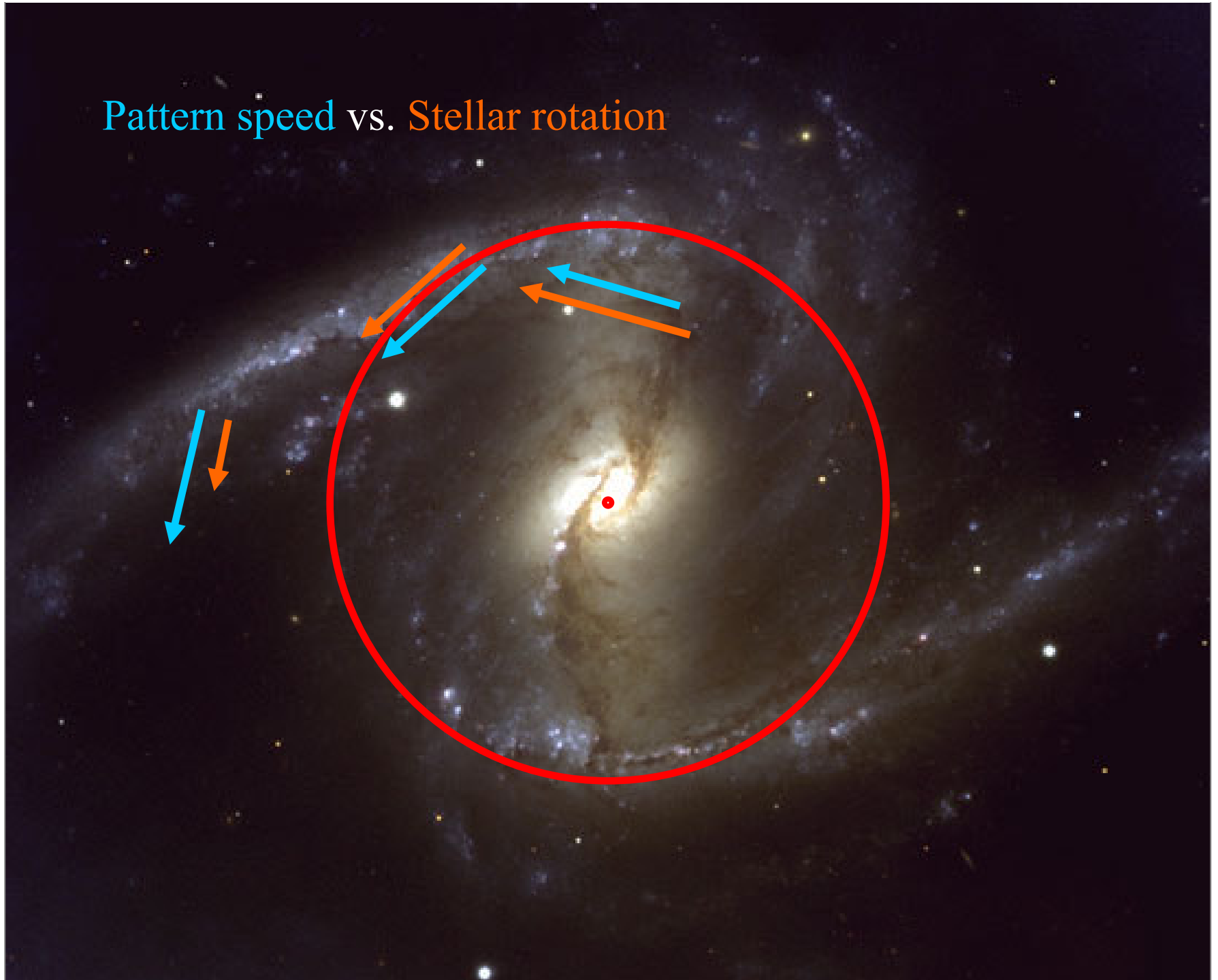
1'st order eq. of state:

$$h_a = c_s^2(\Sigma_0) \frac{\Sigma_a}{\Sigma_0}$$

Pattern speed: $\Omega_p = \frac{\omega}{m}$

Exponent of wave: $i(m\phi - \omega t) = im(\phi - \Omega_p t)$

Pattern speed vs. Stellar rotation



Continuity equation: Gas disk

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

In cylindrical coordinates:

$$\frac{\partial \rho}{\partial t} + \frac{1}{R} \frac{\partial (R \rho v_R)}{\partial R} + \frac{1}{R} \frac{\partial (\rho v_\phi)}{\partial \phi} + \frac{\partial (\rho v_z)}{\partial z} = 0$$

For a disk, 1'st order:

$$\frac{\partial \Sigma_1}{\partial t} + \frac{1}{R} \frac{\partial}{\partial R} (R \Sigma_0 v_{R,1}) + \Omega \frac{\partial \Sigma_1}{\partial \phi} + \frac{\Sigma_0}{R} \frac{\partial v_{\phi,1}}{\partial \phi} = 0$$

For wave solution:

$$i(m\Omega - \omega) \Sigma_a + \frac{1}{R} \frac{d}{dR} (R \Sigma_0 v_{R,a}) + \frac{im\Sigma_0}{R} v_{\phi,a} = 0$$

Summary: 5 equations, 5 unknowns

Euler equation

$$v_{R,a} = -\frac{i}{\Delta} \left[(m\Omega - \omega) \frac{d}{dR} (\Phi_a + h_a) + \frac{2m\Omega}{R} (\Phi_a + h_a) \right]$$

$$v_{\phi,a} = \frac{1}{\Delta} \left[-2B \frac{d}{dR} (\Phi_a + h_a) + \frac{m(m\Omega - \omega)}{R} (\Phi_a + h_a) \right]$$

Continuity equation

$$i(m\Omega - \omega)\Sigma_a + \frac{1}{R} \frac{d}{dR} (R\Sigma_0 v_{R,a}) + \frac{im\Sigma_0}{R} v_{\phi,a} = 0$$

Equation of state

$$h_a = c_s^2(\Sigma_0) \frac{\Sigma_a}{\Sigma_0}$$

+ Poisson Equation

Tight-winding approximation:

$$\Sigma_1(R, \phi, t) = H(R, t)e^{i[m\phi + f(R, t)]}$$

$$\Phi_1(R, \phi, t) = -\frac{2\pi G}{|k|} \Sigma_1(R, \phi, t) \quad (\text{Used Poisson Eq.})$$

$$k(R, t) = \frac{\partial f(R, t)}{\partial R} \quad \cot i = \left| \frac{kR}{m} \right| \gg 1$$

Apply it here:

$$\Sigma_a = H(R)e^{if(R)} \quad k = \frac{df}{dR}$$

$$\Phi_a = -\frac{2\pi G}{|k|} H(R)e^{if(R)}$$

$$h_a = G(R)e^{if(R)} \quad (\text{H, G are smooth, f varies rapidly})$$

Euler equation(s)

$$v_{R,a} = -\frac{i}{\Delta} \left[(m\Omega - \omega) \frac{d}{dR} (\Phi_a + h_a) + \frac{2m\Omega}{R} (\Phi_a + h_a) \right]$$

$$v_{R,a} = \frac{1}{\Delta} (m\Omega - \omega) k (\Phi_a + h_a)$$

$$v_{\phi,a} = \frac{1}{\Delta} \left[-2B \frac{d}{dR} (\Phi_a + h_a) + \frac{m(m\Omega - \omega)}{R} (\Phi_a + h_a) \right]$$

$$v_{\phi,a} = -\frac{2i}{\Delta} B k (\Phi_a + h_a)$$

Continuity equation

$$i(m\Omega - \omega) \Sigma_a + \frac{1}{R} \frac{d}{dR} (R \Sigma_0 v_{R,a}) + \frac{im\Sigma_0}{R} v_{\phi,a} = 0$$

$$(m\Omega - \omega) \Sigma_a + k \Sigma_0 v_{R,a} = 0 \quad [\text{Dropped } 1/(kR) \text{ last term}]$$

5 equations:

$$\left\{ \begin{array}{l} v_{R,a} = \frac{1}{\Delta} (m\Omega - \omega) k (\Phi_a + h_a) \\ v_{\phi,a} = -\frac{2i}{\Delta} B k (\Phi_a + h_a) \\ (m\Omega - \omega) \Sigma_a + k \Sigma_0 v_{R,a} = 0 \\ h_a = c_s^2 (\Sigma_0) \frac{\Sigma_a}{\Sigma_0} \\ \Phi_a = -\frac{2\pi G}{|k|} \Sigma_a \end{array} \right.$$

(Continuity equation)

$$\Rightarrow \Sigma_a \left(1 + \frac{k^2 c_s^2}{\Delta} - \frac{2\pi G \Sigma_0 |k|}{\Delta} \right) = 0$$

$$\Delta \equiv \kappa^2 - (m\Omega - \omega)^2$$

$$\Rightarrow (m\Omega - \omega)^2 = \kappa^2 - 2\pi G \Sigma |k| + k^2 c_s^2$$

Dispersion relation

$$(m\Omega - \omega)^2 = \kappa^2 - 2\pi G\Sigma|k| + k^2 c_s^2 \quad \text{Gas disk}$$

$$(m\Omega - \omega)^2 = \kappa^2 - 2\pi G\Sigma|k| \mathcal{F}\left(\frac{\omega - m\Omega}{\kappa}, \frac{k^2 \sigma_R^2}{\kappa^2}\right)$$

Stellar disk (B&T 6-46)

We have established local solutions ✓

But not a full global solution:

- 1) Assumed tight winding
- 2) Boundary conditions (Disk center & edge)
- 3) Singularities (Resonances)

Basic reason for spiral structure:

Minimize energy while conserving angular momentum

$$(m\Omega - \omega)^2 = \kappa^2 - 2\pi G\Sigma|k| + k^2 c_s^2 \quad \text{Gas disk}$$

Consider now axisymmetric perturbations, $m=0$:

Same equation as long as: $|kR| \gg 1$

$$\omega^2 = \kappa^2 - 2\pi G\Sigma|k| + k^2 c_s^2$$

$$\Sigma \propto e^{-i\omega t}$$

$$Q \equiv \frac{c_s \kappa}{\pi G \Sigma} > 1$$

Stability

$$Q \equiv \frac{\sigma_R \kappa}{3.36 G \Sigma} > 1$$

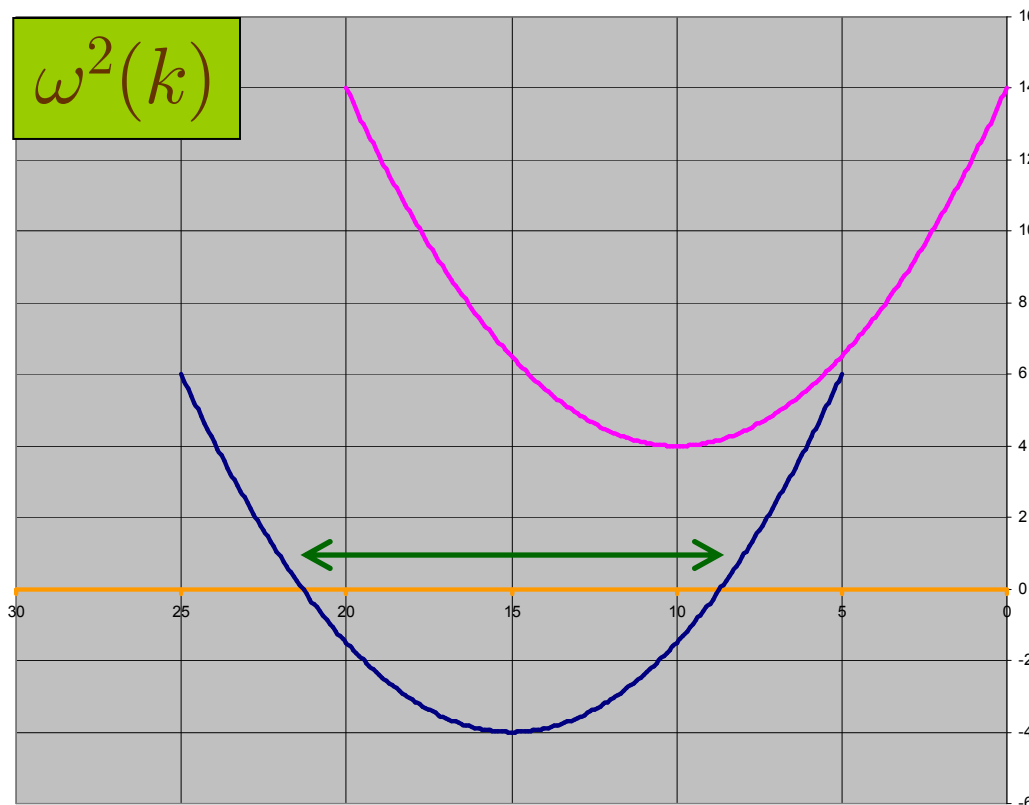
Stellar disk

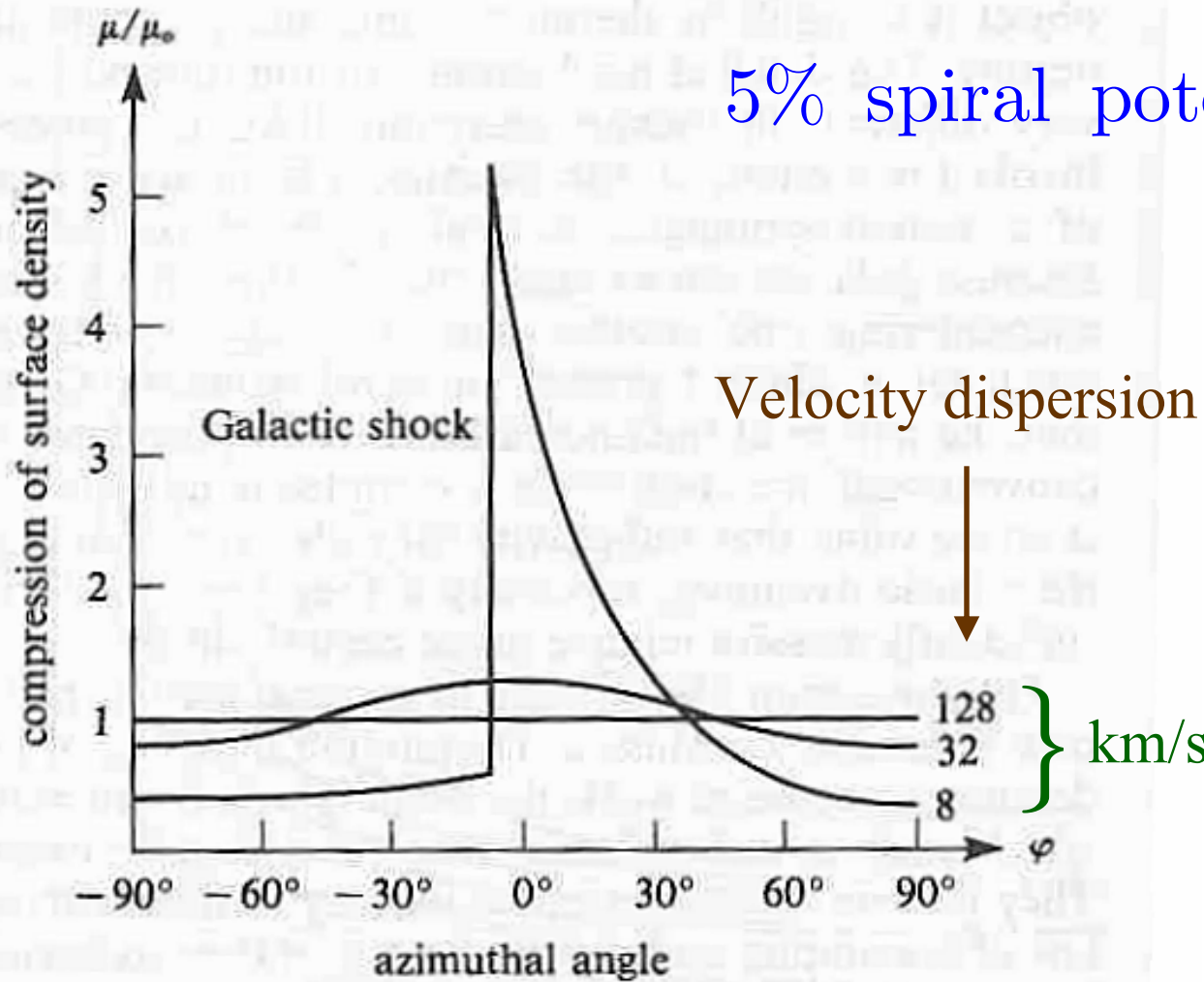
Toomre (1964)

Cooling $\Rightarrow$

Instability $\Rightarrow$

Star Formation





5% spiral potential

Velocity dispersion

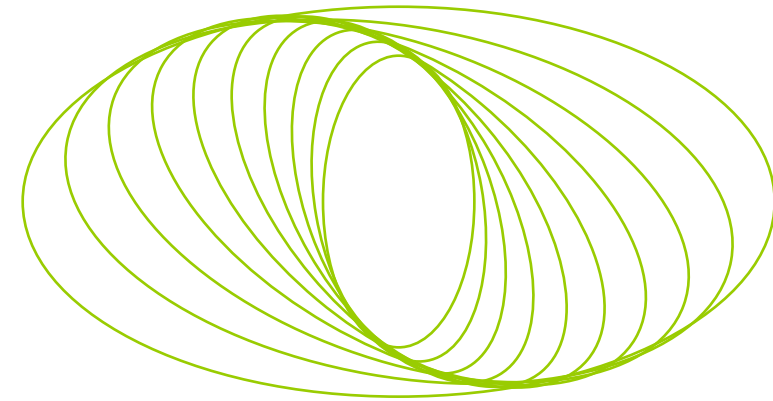
- halo stars
- disk stars
- gas clouds



Spiral Patterns from Elliptical Stellar Orbits

≤ 10 orbits

46 orbits:



NGC2997 AAT 017

