## Cosmology 2 (Prof. Rennan Barkana): Solution to Homework 1

1. In the case without  $\Lambda$ , we can do this analytically (numerically is also OK). Letting  $\Omega_m$  and  $\Omega_r$  be the present  $\Omega$  in matter and radiation, the Friedmann equation is:

$$\left(\frac{1}{a}\frac{da}{dt}\right)^2 = H^2 = H_0^2 \left(\frac{\Omega_m}{a^3} + \frac{\Omega_r}{a^4}\right) ,$$

where  $\Omega_r = 1 - \Omega_m$ . Solving this for dt, we get

$$t(a) = \int dt = H_0^{-1} \int_0^a \frac{da}{\sqrt{\Omega_m/a + \Omega_r/a^2}} = \frac{2}{3\Omega_m^2 H_0} \left[ 2\Omega_r^{3/2} + \sqrt{\Omega_m a + \Omega_r} \left(\Omega_m a - 2\Omega_r\right) \right] .$$

Similarly,

$$\tau(a) = \int \frac{dt}{a} = \frac{2}{\Omega_m H_0} \left( \sqrt{\Omega_m a + \Omega_r} - \sqrt{\Omega_r} \right) \; .$$

With a cosmological constant, we have:

$$t(a) = \int dt = H_0^{-1} \int_0^a \frac{da}{\sqrt{\Omega_m/a + \Omega_r/a^2 + \Omega_\Lambda a^2}}$$

and

$$\tau(a) = \int \frac{dt}{a} = H_0^{-1} \int_0^a \frac{da}{\sqrt{\Omega_m a + \Omega_r + \Omega_\Lambda a^4}}$$

Also note that  $H_0^{-1} = 9.78 \text{ Gyr}/h = 14.4 \text{ Gyr}$ . I will use index 1 for the model without  $\Lambda$ , and 2 with  $\Lambda$ . Today a = 1. In the first model,  $\Omega_m = 1 - 9.07 \times 10^{-5}$ , and in the second model,  $\Omega_m = 1 - 0.689 - 9.07 \times 10^{-5}$ . So matter-radiation equality,  $a = \Omega_r / \Omega_m$  is different in the two models: it is  $a = 9.07 \times 10^{-5}$  in the first model and  $a = 2.92 \times 10^{-4}$  in the second. Other example redshifts: z = 8 (roughly cosmic reionization) and z = 1100 (cosmic recombination). The numerical values (in Gyr units) are  $t_1(9.07 \times 10^{-5}) = 4.87 \times 10^{-6}$ ,  $t_2(2.92 \times 10^{-4}) = 5.04 \times 10^{-5}$ ,  $\tau_1(9.07 \times 10^{-5}) = 0.114$ ,  $\tau_2(2.92 \times 10^{-4}) = 0.367$ ,  $t_1(9.08 \times 10^{-4}) = 2.38 \times 10^{-4}$ ,  $t_2(9.08 \times 10^{-4}) = 3.67 \times 10^{-4}$ ,  $\tau_1(9.08 \times 10^{-4}) = 0.638$ ,  $\tau_2(9.08 \times 10^{-4}) = 0.910$ ,  $t_1(0.111) = 0.356$ ,  $t_2(0.111) = 0.637$ ,  $\tau_1(0.111) = 9.36$ ,  $\tau_2(0.111) = 16.4$ ,  $t_1(1) = 9.63$ ,  $t_2(1) = 13.8$ ,  $\tau_1(1) = 28.6$ ,  $\tau_2(1) = 46.2$ .

2. The smoothed density field is

$$\bar{\delta}(\vec{x}) = \int d^3x_1 W(|\vec{x}_1 - \vec{x}|) \delta(\vec{x}_1) \; .$$

Then

$$\sigma^2 = \left\langle \bar{\delta}(\vec{x}) \,\bar{\delta}(\vec{x}) \right\rangle = \int d^3x_1 \int d^3x_2 \, W(|\vec{x}_1 - \vec{x}|) W(|\vec{x}_2 - \vec{x}|) \left\langle \delta(\vec{x}_1) \delta(\vec{x}_2) \right\rangle$$

Using inverse Fourier transforms and the definition of the power spectrum, we showed in class that

$$\langle \delta(\vec{x}_1)\delta(\vec{x}_2) \rangle = \int \frac{d^3k}{(2\pi)^3} e^{i\vec{k}\cdot(\vec{x}_1-\vec{x}_2)} P(k) \; .$$

We use this in the expression for  $\sigma^2$ , and also we write each term of the form W(r) as the inverse Fourier transform of  $\tilde{W}(k)$ . We get an expression with five integrals:

$$\sigma^{2} = \int d^{3}x_{1} \int d^{3}x_{2} \int \frac{d^{3}k}{(2\pi)^{3}} \int d^{3}k_{1} \int d^{3}k_{2} \, e^{i\vec{k}\cdot(\vec{x}_{1}-\vec{x}_{2})} P(k) \, \frac{\tilde{W}(k_{1})}{(2\pi)^{3}} e^{i\vec{k}_{1}\cdot(\vec{x}_{1}-\vec{x})} \, \frac{\tilde{W}(k_{2})}{(2\pi)^{3}} e^{i\vec{k}_{2}\cdot(\vec{x}_{2}-\vec{x})} \, .$$

The integral over  $\vec{x}_1$  gives  $(2\pi)^3$  times a Dirac delta function of  $\vec{k} + \vec{k}_1$ , and then the  $\vec{k}_1$  integral simply sets  $\vec{k}_1$  equal to  $-\vec{k}$ . We similarly evaluate the  $\vec{x}_2$  and  $\vec{k}_2$  integrals (i.e., we set  $\vec{k}_2$  equal to  $\vec{k}$ ). Thus, we obtain (note that  $\tilde{W}(k)$  only depends on the magnitude of  $\vec{k}$ ):

$$\sigma^2 = \frac{1}{(2\pi)^3} \int d^3k \, \tilde{W}^2(k) P(k) \; .$$

Note that the result does not depend on the starting point  $\vec{x}$  (since this field is statistically homogeneous).

**3.** The variance is:

$$\sigma^2(R) = \int_0^\infty \frac{1}{2\pi^2} k^2 dk P(k) \frac{9}{x^6} [\sin(x) - x\cos(x)]^2 \, dx$$

where x = kR and  $P(k) = Ak/k_{eq}$  for  $k < k_{eq}$ ,  $P(k) = A[k/k_{eq}]^{-3}$  for  $k > k_{eq}$ . Thus,

$$\sigma^{2}(R) = \frac{9}{2\pi^{2}} A k_{\rm eq}^{3} \left\{ \frac{1}{x_{\rm eq}^{4}} \int_{0}^{x_{\rm eq}} \frac{dx}{x^{3}} [\sin(x) - x\cos(x)]^{2} + \int_{x_{\rm eq}}^{\infty} \frac{dx}{x^{7}} [\sin(x) - x\cos(x)]^{2} \right\} ,$$

where  $x_{eq} = k_{eq}R$ .

We calculate that  $k_{eq}8/h = 0.772$ , and then the normalization to  $\sigma(R_0) = 0.81$  gives  $Ak_{eq}^3 = 12.6$ . See the plots below, of the dimensionless quantities  $P(k)k_{eq}^3$  and  $\sigma(R)$ .

